



Orthogonal Polynomials, Moments, Measure Deformation, Dynamical Systems, and SVD Algorithm

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Outline

Basic Dynamics

Splitting and Factorization
Abstraction
Toda Lattice

Measure, Moment and Orthogonality

Orthogonal Polynomials
Measure Deformation

SVD Dynamics

Lotka-Volterra Equation
SVD Solution

Conclusion

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Basic Dynamics

- ▶ Lax dynamics:

$$\begin{aligned}\frac{dX(t)}{dt} &:= [X(t), k_1(X(t))] \\ X(0) &:= X_0.\end{aligned}$$

- ▶ Parameter dynamics:

$$\begin{aligned}\frac{dg_1(t)}{dt} &:= g_1(t)k_1(X(t)) \\ g_1(0) &:= I.\end{aligned}$$

and

$$\begin{aligned}\frac{dg_2(t)}{dt} &:= k_2(X(t))g_2(t) \\ g_2(0) &:= I.\end{aligned}$$

- $k_1(X) + k_2(X) = X$.

Similarity Property

$$X(t) = g_1(t)^{-1} X_0 g_1(t) = g_2(t) X_0 g_2(t)^{-1}.$$

- ▶ Define $Z(t) = g_1(t) X(t) g_1(t)^{-1}$.
- ▶ Check

$$\begin{aligned} \frac{dZ}{dt} &= \frac{dg_1}{dt} X g_1^{-1} + g_1 \frac{dX}{dt} g_1^{-1} + g_1 X \frac{dg_1^{-1}}{dt} \\ &= (g_1 k_1(X)) X g_1^{-1} + g_1 (X k_1(X) - k_1(X) X) g_1^{-1} \\ &\quad + g_1 X (-k_1(X) g_1^{-1}) = 0. \end{aligned}$$

- ▶ Thus $Z(t) = Z(0) = X(0) = X_0$.

Decomposition Property

$$\exp(tX_0) = g_1(t)g_2(t).$$

- ▶ Trivially $\exp(X_0 t)$ satisfies the IVP

$$\frac{dY}{dt} = X_0 Y, Y(0) = I.$$

- ▶ Define $Z(t) = g_1(t)g_2(t)$.
- ▶ Then $Z(0) = I$ and

$$\begin{aligned} \frac{dZ}{dt} &= \frac{dg_1}{dt} g_2 + g_1 \frac{dg_2}{dt} \\ &= (g_1 k_1(X)) g_2 + g_1 (k_2(X) g_2) = g_1 X g_2 \\ &= X_0 Z \quad (\text{by Similarity Property}). \end{aligned}$$

- ▶ By the uniqueness theorem in ODEs, $Z(t) = \exp(X_0 t)$.

Reversal Property

$$\exp(tX(t)) = g_2(t)g_1(t).$$

- ▶ By Decomposition Property,

$$\begin{aligned}g_2(t)g_1(t) &= g_1(t)^{-1} \exp(X_0 t) g_1(t) \\ &= \exp(g_1(t)^{-1} X_0 g_1(t) t) \\ &= \exp(X(t) t).\end{aligned}$$

Abstract QR -type Decomposition

- ▶ Arbitrary subspace decomposition $gl(n) \iff$ Factorization of a *one-parameter semigroup* in the neighborhood of I as the product of two nonsingular matrices , i.e.,

$$\exp(X_0 t) = g_1(t)g_2(t).$$

- Lie algebra decomposition of $gl(n) \iff$ Lie group decomposition of $Gl(n)$ in the neighborhood of I .
- ▶ The product $g_1(t)g_2(t)$ will be called the *abstract $g_1 g_2$ decomposition* of $\exp(X_0 t)$.



Abstract QR-type Algorithm

- ▶ By setting $t = 1$, we have

$$\exp(X(0)) = g_1(1)g_2(1)$$

$$\exp(X(1)) = g_2(1)g_1(1).$$

- ▶ The dynamical system for $X(t)$ is autonomous \implies The above phenomenon will occur at every feasible integer time.
- ▶ Corresponding to the abstract g_1g_2 decomposition, the above iterative process for all feasible integers will be called the *abstract g_1g_2 algorithm*.

Toda Flow

► Lie algebra decomposition:

- $gl(n) = \{\text{skew symmetric}\} \oplus \{\text{upper triangular}\}$.
- Write

$$X = X^o + X^+ + X^-$$

where X^o is the diagonal, X^+ the strictly upper triangular, and X^- the strictly lower triangular part of X .

- Define

$$\Pi_0(X) := X^- - X^{-T}.$$

► The Toda lattice (Symes'82, Deift et al'83):

$$\begin{aligned} \frac{dX}{dt} &= [X, \Pi_0(X)] \\ X(0) &= X_0. \end{aligned}$$



QR Algorithm and Limiting Behavior

- ▶ Sampled at integer times, $\{X(k)\}$ gives the same sequence as does the QR algorithm applied to the matrix $A_0 = \exp(X_0)$.
- ▶ Evolution starting from X_0 ,
 - The flow maintains the spectrum.
 - The construction of the Toda lattice is based on the physics.
 - This is a Hamiltonian system.
 - A certain physical quantities are kept at constant, i.e., this is a *completely integrable* system.
 - Asymptotic behavior can be analyzed via ODE theory.

Orthogonal Polynomials

- ▶ Orthogonal polynomials with respect to a given a measure $\mu(x)$,

$$\int p_k(x)p_\ell(x) d\mu(x) = \delta_{k,\ell}, \quad k, \ell = 0, 1, \dots$$

- ▶ Three-term recurrence relationship.

$$xp_k(x) = a_k p_{k+1}(x) + b_k p_k(x) + a_{k-1} p_{k-1}(x), \quad k = 1, 2, \dots,$$

with $p_{-1}(x) \equiv 0$ and $p_0(x) \equiv 1$.

- ▶ Corresponding monic polynomials $\{\tilde{p}_k(x)\}_k$,

$$x\tilde{p}_k(x) = \tilde{p}_{k+1} + b_k \tilde{p}_k(x) + a_{k-1}^2 \tilde{p}_{k-1}(x).$$

Hankel Determinants

$$H_k := \det \begin{bmatrix} s_0 & s_1 & \dots & s_{k-1} \\ s_1 & s_2 & & s_k \\ \vdots & & & \vdots \\ s_{k-1} & s_k & \dots & s_{2k-2} \end{bmatrix}.$$

- ▶ s_j are moments with respect to μ ,

$$s_j := \int x^j d\mu(x), \quad j = 0, 1, \dots$$



Classical Moment Problem

- $\tilde{p}_k(x)$ is given by (Akhiezer'65, Szegő'75),

$$\begin{aligned} \tilde{p}_k(x) &= \frac{1}{H_k} \det \begin{bmatrix} s_0 & s_1 & \dots & s_k \\ s_1 & s_2 & & s_{k+1} \\ \vdots & & & \vdots \\ s_{k-1} & s_k & \dots & s_{2k-1} \\ 1 & x & \dots & x^k \end{bmatrix}, \\ &= x^k + c_1^{(k)} x^{k-1} + \dots + c_{k-1}^{(k)} x + c_k^{(k)}, \end{aligned}$$



$$c_j^{(k)} = \frac{(-1)^j}{H_k} \det \begin{bmatrix} s_0 & \dots & s_{k-j-1} & s_{k-j+1} & \dots & s_k \\ s_1 & & & & & s_{k+1} \\ \vdots & & \vdots & \vdots & & \vdots \\ s_{k-1} & \dots & s_{2k-j-2} & s_{2k-j} & \dots & s_{2k-1} \end{bmatrix}.$$

Polynomials in Moments

- ▶ By comparing the corresponding coefficients,

$$a_k^2 = \frac{H_k H_{k+2}}{H_{k+1}^2},$$

$$b_k = c_1^{(k)} - c_1^{(k+1)}.$$

- ▶ A classical result.

$$\mu(x) \longleftrightarrow s_j \longleftrightarrow H_k \longleftrightarrow p_k(x).$$

Measure Deformation

- ▶ Rewrite orthogonality in matrix form,

$$\underbrace{\begin{bmatrix} b_0 & a_0 & 0 & & & \\ a_0 & b_1 & a_1 & 0 & & \\ 0 & a_1 & b_2 & a_2 & 0 & \\ & \ddots & \ddots & \ddots & \ddots & \ddots \end{bmatrix}}_J \begin{bmatrix} p_0(x) \\ p_1(x) \\ p_2(x) \\ \vdots \end{bmatrix} = x \begin{bmatrix} p_0(x) \\ p_1(x) \\ p_2(x) \\ \vdots \end{bmatrix}.$$

- ▶ What can be said about

$$\mu(x; t) \longleftrightarrow J(t),$$

if the measure is time dependent?

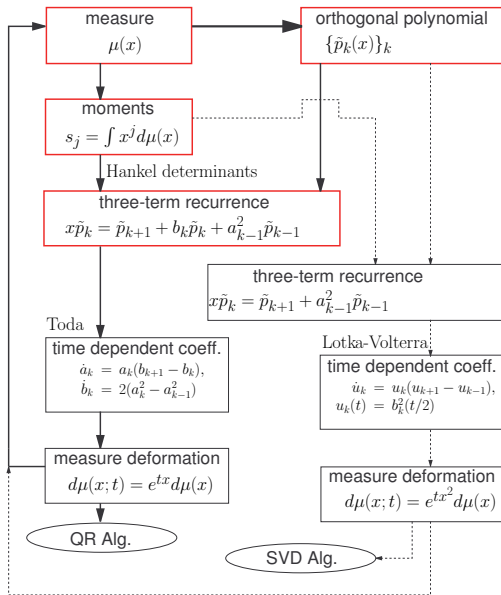
Semi-infinite Toda Lattice

- ▶ If $J(t)$ follows the Toda flow, that is, if

$$\begin{aligned}\frac{da_k}{dt} &= a_k(b_{k+1} - b_k), \\ \frac{db_k}{dt} &= 2(a_k^2 - a_{k-1}^2),\end{aligned}$$

with $a_{-1} \equiv 0$, then (Moser'75)

$$d\mu(x; t) := e^{tx} d\mu(x; 0).$$



Finite-Dimensional Eigenvalue Problem

- ▶ Truncation.

$$\underbrace{\begin{bmatrix} b_0 & a_0 & 0 & & & \\ a_0 & b_1 & a_1 & 0 & & \\ 0 & a_1 & b_2 & a_2 & 0 & \\ & & \ddots & \ddots & \ddots & \\ & & & & & a_{n-2} \\ & & & & 0 & a_{n-2} & b_{n-1} \end{bmatrix}}_L \begin{bmatrix} p_0(x) \\ p_1(x) \\ p_2(x) \\ \vdots \\ p_{n-1}(x) \end{bmatrix} + \begin{bmatrix} 0 \\ 0 \\ 0 \\ \vdots \\ a_{n-1} p_n(x) \end{bmatrix} = x \begin{bmatrix} p_0(x) \\ p_1(x) \\ p_2(x) \\ \vdots \\ p_{n-1}(x) \end{bmatrix}.$$

- ▶ λ is a root of the polynomial $p_n(x)$ if and only if λ is an eigenvalue of the finite-dimensional tridiagonal matrix L .

SVD Flow

- ▶ Given a bidiagonal matrix B_0 ,

$$\frac{dB}{dt} = B\Pi_0(B^\top B) - \Pi_0(BB^\top)B, \quad B(0) = B_0, \quad (1)$$

- ▶ $B(t)$ stays bidiagonal for all t .
- ▶ The sequence $\{B(\ell)\}_\ell$ by sampling $B(t)$ at integer times corresponds to the iterates produced by the Golub-Kahan SVD algorithm.

Lotka-Volterra Equation

- ▶ Denote

$$B(t) := \text{diag} \left\{ \begin{array}{ccccccc} & & b_2(t) & & & & \\ & & & & \dots & & \\ & b_1(t) & & & & & \\ & & & b_3(t) & & & \\ & & & & \dots & & \\ & & & & & & b_{2n-2}(t) \\ & & & & & & & \\ & & & & & & & b_{2n-1}(t) \end{array} \right\},$$

- ▶ Change variables,

$$u_{2k-1}(t) := b_{2k-1}^2 \left(\frac{t}{2} \right),$$

$$u_{2k}(t) := b_{2k}^2 \left(\frac{t}{2} \right).$$

- ▶ *Continuous-time finite Lotka-Volterra equation,*

$$\frac{du_k}{dt} = u_k(u_{k+1} - u_{k-1}), \quad k = 1, 2, \dots, 2n-1, \quad (2)$$

with $u_0(t) \equiv 0$ and $u_{2n}(t) \equiv 0$.



τ Functions

- ▶ Change variables,

$$u_k = \frac{\tau_{k+2}\tau_{k-1}}{\tau_{k+1}\tau_k}.$$

- ▶ A compatibility condition,

$$\frac{d\tau_k}{dt}\tau_{k+1} - \tau_k \frac{d\tau_{k+1}}{dt} + \tau_{k-1}\tau_{k+2} = 0.$$

Determinantal Solution

- Starting with $\tau_{-1} \equiv 0$, $\tau_0 \equiv 1$, $\tau_1(t) = 1$ and $\tau_2(t) = \psi(t)$,

$$\tau_3 = \frac{d\psi}{dt},$$

$$\tau_4 = \det \begin{bmatrix} \psi & \psi^{(1)} \\ \psi^{(1)} & \psi^{(2)} \end{bmatrix},$$

- In general (Tsujiimoto'95),

$$\tau_{2k-1} = \bar{H}_{k-1,1},$$

$$\tau_{2k} = \bar{H}_{k,0},$$

where

$$\bar{H}_{k,j}(t) := \det \begin{bmatrix} \psi^{(j)} & \psi^{(j+1)} & \dots & \psi^{(j+k-1)} \\ \psi^{(j+1)} & \psi^{(j+2)} & \dots & \psi^{(j+k)} \\ \vdots & \vdots & & \vdots \\ \psi^{(j+k-1)} & \psi^{(j+k)} & & \psi^{(j+2k-2)} \end{bmatrix}, \quad j = 0 \text{ or } 1,$$

is the determinant of a $k \times k$ Hankel matrix.

SVD Solution

- ▶ The general solution to the Lotka-Volterra equation (Tsujiimoto, Nakamura&Iwasaki'01)

$$u_{2k-1}(t) = \frac{\bar{H}_{k,1}(t)\bar{H}_{k-1,0}(t)}{\bar{H}_{k,0}(t)\bar{H}_{k-1,1}(t)},$$

$$u_{2k}(t) = \frac{\bar{H}_{k+1,0}(t)\bar{H}_{k-1,1}(t)}{\bar{H}_{k,1}(t)\bar{H}_{k,0}(t)}, \quad k = 1, 2, \dots, n,$$

- ▶ Assuming all derivatives of ψ are obtainable from elementary calculus,
 - In principle, all Hankel determinants can be calculated algebraically.
 - All quantities about $u_k(t)$ are now in the analytic form.
 - The SVD flow and, hence, the iterates from the SVD algorithm are representable in closed form.



Conclusion

- ▶ Each of the Toda lattice and the Lotka-Volterra equation governs the evolution of a certain class of orthogonal polynomials whose orthogonality is determined by a specific time-dependent measure.
- ▶ Since the measure deformation is explicitly known, moments can be calculated which, when properly assembled, lead to the conclusion abstractly, but literally, that the iterates of the QR algorithm and the SVD algorithm can be expressed in closed-form!
- ▶ Hankel determinantal solutions are too complicated to be useful. However, a “smart” integrability-preserving discretization of the Lotka-Volterra equation can yield a new SVD algorithm.

A New Chapter

- ▶ A key step in the integrable discretization of the Lotka-Volterra equation (2) is a particular Euler-type scheme of the form (Hirota, Tsujimoto&Imai'93),

$$u_k^{[\ell+1]} = u_k^{[\ell]} + \delta \left(u_k^{[\ell]} u_{k+1}^{[\ell]} - u_k^{[\ell+1]} u_{k-1}^{[\ell+1]} \right).$$

- $u_k^{[\ell]} \approx u_k(\ell\delta)$.
- Boundary conditions $u_0^{[\ell]} \equiv 0$ and $u_{2n}^{[\ell]} \equiv 0$ for all ℓ .